A New Nickel(III) Oxide Family: MSr_3NiO_6 (M = Sc, In, Tm, Yb and Lu)

Michael James and J. Paul Attfield*

Abstract: The new phases $M^{III}Sr_3Ni^{III}O_6$ have been prepared for M = Sc, In, Tm, Yb and Lu. Thermogravimetric analysis indicates that these phases are stoichiometric nickel(III) oxides. Rietveld refinement of their crystal structures from powder X-ray diffraction data confirms that they adopt the rhombohedral K_4CdCl_6 -type structure (space group $R\overline{3}c$, a = 9.6595(2) and c = 10.8546(3) Å for $ScSr_3NiO_6$). The M site is fully occupied for M = Sc and In, but a deficiency of

scattering for M = Tm, Yb and Lu is shown to be due to Ni substitution through a simultaneous refinement of the YbSr₃NiO₆ structure using X-ray and

Keywords

Jahn-Teller distortions · magnetic properties · neutron powder diffraction · nickel oxides · X-ray powder diffraction

time-of-flight neutron diffraction data. The refined composition is $(Yb_{0.83}Ni_{0.17})$ - Sr_3NiO_6 . The magnetic susceptibilities of the M=Sc, In and Lu samples show Curie–Weiss behaviour down to 6~K; however, $ScSr_3NiO_6$ shows a broad transition between 250 and 290 K, with Curie–Weiss behaviour above and below this anomaly. This transition is thought to be between the statically and dynamically Jahn–Teller distorted regimes of octahedrally coordinated, low-spin Ni^{3+} .

Introduction

Nickel(III) oxides are of interest for their "high" oxidation state and for their structural, magnetic and electronic properties in relation to high Tc superconductors, as both low spin Ni³ + and Cu² + have $S = {}^1/{}_2$ ground states. The three-dimensional perovskites LnNiO₃ (Ln = La-Gd)[1-3] and layered K₂NiF₄-type Ln₂- ${}_x$ Sr ${}_x$ NiO₄- ${}_\delta$ solid solutions (Ln = La,[4-6] Nd,[7-9] Pr, Sm and Gd)[9,10] have been extensively studied. We have recently prepared a series of related NiIII defect oxides Ln₀₃₃Sr₁,6↑NiO₃,6↑ (LnSr₅Ni₃O₁₁) for Ln = Y, Dy, Ho, Er and Tm.[11,12] Attempts to prepare the Yb analogue of these phases resulted in the formation of a new rhombohedral NiIII compound YbSr₃NiO₆, for which preliminary results have been reported.[13] This paper describes further results for YbSr₃NiO₆ and the preparation and characterisation of the family of isostructural phases MSr₃NiO₆ (M = Sc, In, Tm and Lu).

Results

Powder X-ray diffraction profiles were recorded for each of the new phases. It was found that the major phase in each case for M = Sc, In, Tm, Yb and Lu could be indexed on a rhombohe-

dral unit cell with hexagonal lattice parameters similar to those of K_4CdCl_6 -type oxides such as $Sr_4PtO_6.^{[14,\,15]}\,Small$ quantities of M_2SrO_4 phases $^{[16-18]}$ were observed for $M=Tm,\,Yb$ and Lu, and SrO was found in all of the samples containing the rhombohedral phase, with the exception of $InSr_3NiO_6$, which was phase-pure by powder X-ray diffraction. The indexed pattern of $InSr_3NiO_6$ is given in Table 1. The rhombohedral phase was not observed for M=Ga, which gave a mixture of phases including $Sr_5Ni_4O_{11},^{[19]}\,Sr_3Ga_2O_6,^{[20]}\,NiO$ and SrO, or for M=Er, which gave $ErSr_5Ni_3O_{11},^{[12]}\,Er_2SrO_4,^{[16]}$ and SrO.

Table 1. Observed d spacings and relative peak intensities (I_t) for X-ray diffraction pattern of ${\rm InSr_3NiO_6}$.

hkl	d spacing (Å) I _r	hkl	d spacing (Å)	$I_{\rm r}$
110	4.787	15	223	2.014	23
102	4.620	1	321	1.883	2
202	3.320	13	006	1.847	7
211	3.022	4	410	1.842	4
113	2.923	48	215	1.818	4
300	2.772	100	322	1.807	11
212	2.733	8	314	1.773	3
104	2.624	3	116	1.722	1
204	2.303	19	404	1.667	1
311	2.260	17	413	1.630	5
312	2.134	1	330	1.604	11
214	2.080	17			

The thermogravimetric reduction plot of InSr₃NiO₆ is shown in Figure 1. The mass loss upon reduction of this phase was 8.9(1)%, corresponding to an oxygen content of 5.97(1) per formula unit. The small plateau at around 1.7% mass loss suggests that Ni³⁺ is initially reduced to Ni²⁺, before In³⁺ and Ni²⁺ are together reduced to their metallic states.

Department of Chemistry, University of Cambridge Lensfield Road, Cambridge CB21EW (UK) and

Interdisciplinary Research Centre in Superconductivity University of Cambridge, Madingley Road, Cambridge CB30HE (UK) Fax: Int. code +(1223)336-362 e-mail: jpa14@cam.ac.uk

Dr. M. James

Neutron Scattering Group, ANSTO, Lucas Heights Research Laboratories Private Mail Bag 1, Menai, N. S. W. 2234 (Australia)

^[*] Dr. J. P. Attfield

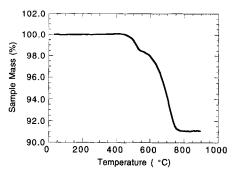


Fig. 1. Thermogravimetric reduction profile for InSr₃NiO₆.

Rietveld refinements of the MSr₃NiO₆ structures were carried out on powder X-ray diffraction data. The starting model in space group $R\bar{3}c$ was that previously used for YbSr₃NiO₆ [13] and is derived from the structure of Sr₄PtO₆^[14, 15] with M at the unique trigonal prismatic Sr site and Ni at the Pt site (Fig. 2). The impurity phases were also fitted by refining their scale factors and lattice parameters. Variation of the site occupancies

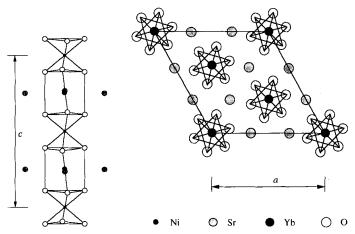


Fig. 2. Left: NiYbO₆ chains. Right: unit cell for YbSr₃NiO₆

showed that all the sites are fully occupied in ScSr₃NiO₆ and $InSr_3NiO_6$, but for MSr_3NiO_6 with M = Tm, Yb and Lu a deficiency of scattering of about 12% was observed at the M site. In these cases, the observed electron density at the M site is equally consistent with ca. 12% vacancy formation, ca. 18% Ni substitution or ca. 26% Sr substitution. The latter possibility can be discarded as the average "M"-O distance is less then that expected for M alone, whereas Sr substitution would expand this site.

To distinguish between the vacancy formation and Ni substitution models, the structure of YbSr₃NiO₆ was refined by using time-of-flight neutron and X-ray powder diffraction profiles simultaneously. In the latter stages of this refinement, Yb and Ni atoms were placed at the M site and their occupation factors were freely refined, giving occupation factors of 0.84(1) and 0.15(1), respectively. The total Yb+Ni occupancy of 0.99(2)does not differ significantly from unity, showing that Ni substitution and not vacancy formation occurs. The total Yb+Ni occupancy was thus constrained to be 1.0 in the final model. There was no evidence for disorder or vacancies at any other

The results of the powder X-ray refinements are given in Table 2, and X-ray diffraction profiles of InSr₃NiO₆ are shown in Figure 3. The derived variations of c/a and V with M^{3+} ionic

Table 2. Profile and structural parameters, interatomic distances and angles, and impurity phase ratios for MSr₃NiO₆ (M = Sc, In, Lu, Yb and Tm) (e.s.d.'s in parenthe-

	Sc	In	Lu	Yb	Tm
Cell paramete	rs				-
a (Å)	9.6595(2)	9.6458(3)	9.6856(2)	9.6851(2)	9.6838(2)
c (Å)	10.8546(3)	11.1111(3)	11.0145(2)	11.0960(2)	11.1193(3)
$V(\mathring{A}^3)$	877.10(4)	895.29(6)	899.57(4)	901.37(4)	903.02(5)
R factors (%)	ı				
R_{WP}	4.8	6.5	5.4	2.8	6.7
$R_{\rm p}$	2.9	4.2	3.5	1.9	4.4
$R_{\rm F}$	4.1	3.5	2.9	4.1	5.5
Atomic paran	neters [a]				
Ni $U_{\rm iso}$ (Å ²)	0.006(1)	0.002(2)	0.005(1)	0.005(1)	0.008(2)
Sr x	0.3709(1)	0.3705(1)	0.3701(1)	0.3704(1)	0.3698(2)
Sr U_{iso} (Å ²)	0.006(1)	0.007(1)	0.006(1)	0.005(1)	0.008(1)
$M U_{iso} (Å^2)$	0.001(1)	0.005(1)	0.001(1)	0.005(1)	0.002(1)
M/Ni occup.	1.00	1.00	0.81(1)/0.19	0.83(1)/0.17	0.81(1)/0.19
O x	0.1746(6)	0.1781(6)	0.1759(6)	0.1785(6)	0.1755(8)
Oy	0.0220(5)	0.0232(8)	0.0228(7)	0.0249(8)	0.0238(9)
O z	0.1158(3)	0.1140(6)	0.1112(5)	0.1113(6)	0.1108(7)
O U _{iso} (Å ²)	0.002(1)	0.001(2)	0.005(2)	0.008(2)	0.009(2)
Interatomic d	istances (Å)				
$Ni-O \times 6$	2.028(5)	2.055(6)	2.022(5)	2.039(6)	2.017(7)
$M-O\times 6$	2.157(5)	2.214(6)	2.222(5)	2.235(6)	2.224(7)
$Sr-O \times 2$	2.483(5)	2.489(6)	2.523(5)	2.516(6)	2.534(7)
$Sr-O \times 2$	2.610(5)	2.629(6)	2.603(5)	2.599(6)	2.612(8)
$Sr-O\times 2$	2.649(5)	2.637(6)	2.644(6)	2.629(6)	2.634(8)
$Sr-O \times 2$	2.690(5)	2.696(6)	2.703(6)	2.718(6)	2.714(8)
mean	2.608(5)	2.613(6)	2.618(6)	2.616(6)	2624(8)
Interatomic a	ngles (°) [b]				
O-Ni-O'	85.6(2)	86.0(3)	86.8(2)	87.1(3)	86.6(3)
O-Ni-O"	94.4(2)	94.0(3)	93.2(2)	92.9(3)	93.4(3)

[a] See Table 3 for atomic positions. [b] z(O) = z(O') = -z(O'').

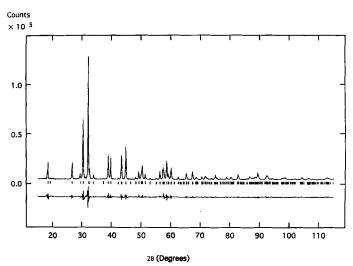


Fig. 3. Observed (points), calculated (full line) and difference X-ray powder diffraction patterns for InSr₃NiO₆

radius are shown in Figure 4, and the twist angles (ϕ) for the distorted MO₆ trigonal prisms (defined in Fig. 5, left) are plotted in Figure 5 (right). The results of the simultaneous refinement of YbSr₃NiO₆ are shown in Table 3, and the observed, calculated and difference neutron diffraction profiles are shown in Figure 6. The associated X-ray diffraction profiles are essentially identical to those shown in reference [13].

The molar susceptibility of YbSr₃NiO₆ has been previously reported. ^[13] The magnetisation of MSr_3NiO_6 samples (M = Sc,

Table 3. Profile and structural parameters from the simultaneous refinement of $Yb_{1-x}Ni_xSr_3NiO_6$ from X-ray and TOF neutron diffraction data in space group R3c with e.s.d.'s in parentheses (cell dimensions (Å): a = 9.6831(3), c = 11.0913(3)).

a) R factors (%).

			R_{wp}	R	P	$R_{ m F}$
X-ray TOF neutron		3.1 2.0 2.1 3.6			4.2 1.4	
				o ————		
b) Aton	nic para	meters.				
Atom	Posi- tion	x	у	z	$U_{\mathrm{iso}}\ (\mathrm{\mathring{A}}^2)$	Site occup.
Ni	6 <i>b</i>	0	0	0	0.0056(2)	1.0
Yb/Ni	6 <i>a</i>	0	0	0.25	0.0051(2)	0.833(8)/0.167
Sr	18 <i>e</i>	0.3707(1)	0	0.25	0.0073(1)	1.0
O	36 <i>f</i>	0.1748(1)	0.0226(1)	0.1106(1)	0.0110(2)	1.0
c) Intera	atomic (distances (Å)	and angles	(°)		
Ni-O	×	6	2.0112(5)	Yb/Ni-	O ×6	2.2208(4)
Sr-O	×	2	2.5408(6)	Sr-O	×2	2.6048(5)
Sr-O	×	2 .	2.6437(5)	Sr-O	$\times 2$	2.6971(7)
O-Ni-O'	×	6 8	6.68(2)	O-Ni-O"	× 6	93.32(2)
O'-Ni-O	" ×	3 18	0.0			

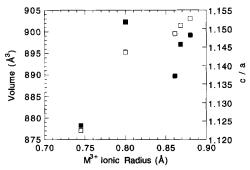


Fig. 4. The variation of cell volume V (open squares) and c/a (filled squares) with M^{3+} ionic radius for MSr_3NiO_6 (M=Sc, In, Lu, Yb and Tm in order of increasing radius).

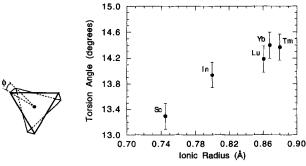


Fig. 5. Left: Schematic diagram of distorted MO_6 trigonal prism defining the twist angle (ϕ) . Right: Plot of twist angle versus M^{3+} ionic radius for MSr_3NiO_6 .

In and Lu) were initially measured between 6 and 325 K. The inverse molar susceptibilities $(1/\chi_{\rm M})$, calculated from these data, are shown in Figure 7. ${\rm InSr_3NiO_6}$ and ${\rm LuSr_3NiO_6}$ show Curie–Weiss behaviour throughout, whereas ${\rm ScSr_3NiO_6}$ displays a broad magnetic transition between around 250 and 290 K. To clarify the magnetic behaviour of ${\rm ScSr_3NiO_6}$ above the anomaly, the susceptibility of a second sample encapsulated by a quartz holder was measured up to 400 K. Effective magnetic moments ($\mu_{\rm eff}$) and Weiss parameters (θ) for the M = In, Yb and Lu phases, and for ${\rm ScSr_3NiO_6}$ below 200 K and above 300 K are given in Table 4.

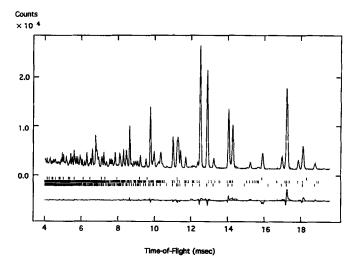


Fig. 6. Observed (points), calculated (full line) and difference time-of-flight neutron powder diffraction patterns for $YbSr_3NiO_6$. Reflection bars for the principal phase and the impurities Yb_2SrO_4 and SrO are shown.

Table 4. Effective moments (μ_{eff}) and Weiss parameters (θ) for MSr₃NiO₆ (M = Yb, Lu, In and Sc) fitted over the given temperature ranges.

M	T(K)	$\mu_{\rm eff}$ (B. M.)	θ (K)	
Yb	150-300	4.6(1)	-29(1)	
Lu	6 - 300	2.30(1)	3.4(1)	
In	6-300	1.97(1)	2.4(1)	
Sc (sample 1)	6 - 200	2.20(1)	3.2(1)	
Sc (sample 2)	6-200	2.18(1)	4.4(1)	
Sc (sample 2)	300-400	2.11(1)	1.8(1)	

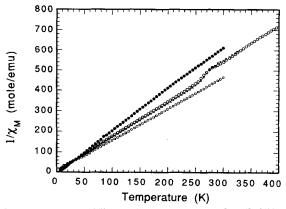


Fig. 7. The inverse molar susceptibility $(1/\chi_M)$ versus temperature for MSr_3NiO_6 where M=Sc (two samples; crosses and open circles), In (filled circles) and Lu (open squares).

Discussion

Under the preparative conditions used, rhombohedral $K_4\text{CdCl}_6$ -type phases $MSr_3\text{NiO}_6$ are stabilised by a range of M (= Sc, In, Lu, Yb and Tm) at the trigonal prismatic site with ionic radii^[21] $r^{(\text{VI}M^{3+})}$ between 0.75 and 0.87 Å. Attempts to form these types of phases under the same conditions with $M = Ga^{3+}$ and Er^{3+} (having $r^{(\text{VI}M^{3+})} = 0.62$ and 0.89 Å, respectively), were unsuccessful. The kinetics of formation of these phases also seem to depend upon $r^{(\text{VI}M^{3+})}$. The smaller In^{3+} and Sc^{3+} reacted to completion in 20–30 hours, but the larger lanthanide cations took in excess of 150 hours to react with many intermediate regrindings.

Rietveld refinement using X-ray powder diffraction data confirms that MSr_3NiO_6 (M = Sc, In, Lu, and Tm) are isostructural with $YbSr_3NiO_6$ (Fig. 2). In this phase $NiYbO_6$ chains are formed down the threefold z axis by face-sharing between alternating NiO_6 octahedra and slightly twisted YbO_6 trigonal prisms (Fig. 2, left). Strontium, with an irregular eightfold coordination, occupies positions between the $NiYbO_6$ chains (Fig. 2, right). The unit cell volume increases with $r(^{VI}M^{3+})$ (Fig. 4). However, for $InSr_3NiO_6$, a is smaller and c is larger than expected; this is clearly seen in the accompanying plot of c/a.

The simultaneous X-ray and neutron study of YbSr₃NiO₆ confirms the oxygen stoichiometry. The ratio of X-ray $(f_{Yb}(0):f_{Ni}(0)=2.50)$ and neutron $(b_{Yb}:b_{Ni}=1.22)$ scattering factors for Yb and Ni are sufficiently different that good contrast was obtained in the simultaneous refinement, enabling the substitution of Ni for Yb to be demonstrated. The M = Tm, Yb and Lu phases are of composition $(M_{1-x}Ni_x)Sr_3NiO_6$ with $x\approx0.17$, but for M = Sc or In x is 0.

The rhombohedrally distorted NiO₆ octahedra (D_{3d} symmetry) in the MSr₃NiO₆ structures are elongated parallel to c, and the Ni–O bond lengths and angles (Table 2) vary little with M. A similar Ni³⁺ environment is found in LiNiO₂, ^[122] with six equivalent Ni–O distances of 1.969 (1) Å, and cis O-Ni-O bond angles of 86.1 and 93.9°. The increase in M–O distance on going from M = Sc to Tm is less than would be expected on the basis of their ionic radii. For the larger cations the mean $r(^{VI}M^{3+})$ is reduced by approximately 6% through the substitution of Ni, and by increasing the twist angle of the trigonal pyramid (Fig. 5, right), which enables the triangular faces to approach each other more closely. Less size pressure at the M site for M = Sc and In does not result in any observable Ni substitution.

These compounds are the first $M^{III}Sr_3T^{III}O_6$ oxides (T = transition metal) to be synthesised with the rhombohedral $R\bar{3}c$ symmetry of a K_4CdCl_6 -type^[23] structure. Previously only $M^{II}Sr_3T^{IV}O_6$ oxides such Sr_4TO_6 (T = Pt, Ir and Rh), II4] Ca_4TO_6 (T = Pt, I²⁴] and $Ir^{[25]}$) and $Ni^{II}Sr_3Pt^{IV}O_6$ [26] have been reported in which T^{IV} occupies the octahedral sites and Sr^{II} , Ca^{II} or Ni^{II} , are present at the trigonal prismatic sites. It is notable that this structure type enables alkaline earth (Sr, Ca), transition metal (Sc, Ni), lanthanide (Tm, Yb, Lu) and main-group (In) cations to be observed in the unusual trigonal-prismatic coordination geometry. An isostructural disordered phase $Cu_{0.75}Ca_{0.25}(Sr_{1.5}Ca_{1.5})PtO_6$ has also been reported, I^{27]} as well as $CuSr_3PtO_6$ and $CuSr_3IrO_6$, which crystallise in the monoclinic subgroup C2/c. In these structures, Cu^{II} is displaced from the centre of the trigonal prism towards one of the square faces; this results in a typical CuO_4 square-planar geometry.

Ternary M^{II} –T^{IV} oxides can display the perovskite-type MTO₃, K₂NiF₄-type M₂TO₄ and K₄CdCl₆-type M₄TO₆ structures as the M/T ratio increases. The relationships between these and other hexagonal perovskite structures have recently been described.^[29] The same progression of structures may now be generated by the double substitution of Ln^{III} and Ni^{III} for M^{II} and T^{IV} giving LnNiO₃, LnSrNiO₄ and LnSr₃NiO₆. The former two structures are stabilised by large lanthanide cations (Ln = La-Gd), but the latter is stabilised by the cation ordering of small Ln (= Tm-Lu) and Sr cations.

The MSr_3NiO_6 structure provides a good opportunity to study the magnetic properties of Ni^{3+} within isolated NiO_6 octahedra. We have previously shown that when paramagnetic Yb^{3+} is present at the M site a broad maximum is observed at around 12 K, suggesting possible one-dimensional magnetic ordering due to antiferromagnetic superexchange interactions within the $YbNiO_6$ chains. This is supported by the Weiss constant $\theta = -29$ K extracted from the high-temperature re-

gion where the magnetic susceptibility follows the Curie–Weiss law. With diamagnetic Sc, In or Lu present in the trigonal prismatic sites between the NiO₆ octahedra, there is no evidence for any short-range magnetic order down to 4 K (Fig. 7), and small values of θ in the range 2–4 K are found. The $\mu_{\rm eff}$ values in Table 4 are all characteristic of low-spin 3 d⁷ Ni³⁺; however, the broad transition observed for ScSr₃NiO₆ between 250 and 290 K suggests that two subtly different limiting forms of magnetic behaviour are possible for these phases. For LuSr₃NiO₆ and ScSr₃NiO₆ (<250 K), $\mu_{\rm eff}$ and θ have slightly greater values than for InSr₃NiO₆ and ScSr₃NiO₆ (>290 K). The anomalous c/a ratio for InSr₃NiO₆ (Fig. 4) in comparison to those those for the M = Lu, Yb and Tb analogues may also reflect a subtle electronic difference in the behaviour of Ni³⁺ in these materials.

The rhombohedral distortion of the NiO₆ octahedra in the MSr_3NiO_6 structure does not remove the degeneracy of the 2E_a ground state of Ni³⁺. This state is prone to a Jahn-Teller distortion, and so the two magnetic behaviours may be identified with statically (<250 K) and dynamically (>290 K) distorted NiO₆ octahedra. In the former case, the small distortions are not ordered within the structure, so that no change of crystal structure is seen. The proposed dynamic-to-static transition in Sc-Sr₃NiO₆ is similar to that reported in K₂PbCu(NO₂)₆. The Cu²⁺ ESR spectrum of this compound shows a gradual transition over the range 293-273 K. [30, 31] The room-temperature spectrum gave an isotropic g value of 2.10, but at liquid-nitrogen temperatures g was clearly anisotropic with estimated values of $g_{\perp} = 2.155$ and $g_{\parallel} = 2.061$, suggesting a freezing-out of the dynamic distortion to give a static one. The equivalent isotropic g value of 2.12 at 80 K is slightly greater than that at 300 K, consistent with the slight increase in μ_{eff} observed for ScSr₃NiO₆ on cooling through the transition. Variable-temperature neutron diffraction and ESR studies of ScSr₃NiO₆ are under way to clarify the nature of this unusual transition.

Conclusions

This study has shown that rhombohedral, K_4CdCl_6 -type, nickel(III) oxides MSr_3NiO_6 are formed for M=Sc, In, Tm, Yb and Lu at $1100\,^{\circ}C$. These are the first oxide phases to adopt this structure type with trivalent ions on both the trigonal-prismatic and octahedral sites. Simultaneous refinement of the structure of $YbSr_3NiO_6$ using X-ray and time-of-flight neutron diffraction data shows that approximately 17% Ni is substituted for Yb at the M site. This substitution is also found for M=Tm and Lu and helps to stabilise the structure by reducing the average ionic radius at the M site.

The magnetic susceptibility curve of YbSr₃NiO₆ shows evidence of one-dimensional magnetic interactions within the NiYbO₆ chains, but this is not seen with diamagnetic M=Sc, In or Lu. The magnetic susceptibility of $ScSr_3NiO_6$ shows an anomaly between 250 and 290 K, with slightly different Curie—Weiss limiting behaviours above and below this temperature region. This transition is consistent with a change from dynamic to static Jahn–Teller distortions of the Ni^{II}O₆ octahedra on cooling.

Experimental Procedure

Sample preparation: Polycrystalline samples with bulk composition MSr_3NiO_6 (M = Sc, In, Ga, Lu, Yb, Tm and Er) were synthesised from spectroscopic-grade powders of strontium carbonate, nickel nitrate hexahydrate and the corresponding metal oxide M_2O_3 (M = Lu, Yb, Tm and Er) or nitrate pentahydrate $M(NO_3)_3 \cdot 5H_2O$ (M = Sc, In and Ga). Prior to weighing the lanthanide oxides were preheated to $1000\,^{\circ}C$ in air to decompose any carbonate material to the oxide. The

powders were dissolved in dilute nitric acid, and an intimate mixture of the metal oxides was formed by decomposition of a citric acid/ethylene glycol gel. The residues were pelleted and sintered in a tube furnace at $1100\,^{\circ}\text{C}$ under flowing oxygen for up to one week with frequent regrinding and repelleting until no further reaction was evident by powder X-ray diffraction. The reaction times for M = Sc and In (20–30 h) were substantially lower than for M = Tm-Lu (ca. 150 h).

Powder Diffraction: Powder X-ray diffraction profiles were recorded on a Philips PW 1710 diffractometer with $Cu_{K\alpha}$ radiation. Data of sufficient quality for structure refinement were collected over $13 \le 2\theta \le 113^\circ$, in 0.025° steps, with integration times of 12 s. These structural refinements were carried out by the Rietveld method [32] using the GSAS program [33] and a refined background function.

The time-of-flight (TOF) powder neutron diffraction profile of YbSr $_3$ NiO $_6$ was collected in 3 h on the POLARIS instrument at the Rutherford Appleton Laboratory. The backscattering profile (135 \leq 2 θ \leq 158°) was used in a simultaneous Rietveld refinement of the structure with the above X-ray diffraction data. The diffraction pattern of a Si powder standard was used to calibrate the time-of-flight scale.

Thermogravimetric Analysis: Thermogravimetric analysis of a ≈ 30 mg sample of $InSr_3NiO_6$ was carried out with a Stanton Redcroft STA 1500 simultaneous thermal analyser. The sample was reduced under a 5% hydrogen in nitrogen mixture (flow rate of 58 mL min⁻¹) over a temperature range of 15 to 900 °C at a heating rate of 10 °C min⁻¹.

Magnetic Susceptibility Measurements: Magnetic susceptibilities were measured using a Quantum Design SQUID magnetometer under an applied field of 3.0 T. Samples were cooled down to 6 K in zero field, and the magnetisation was measured while they were warmed up to 400 K.

Acknowledgements: The authors wish to thank Dr. Dolores Marcos and Jon Chapman for their assistance in collecting the time-of-flight powder neutron diffraction pattern of YbSr₃NiO₆. M. J. wishes to thank the University of Sydney and Newman College, Melbourne University, for their support.

Received: October 19, 1995 [F 229]

- G. Demazeau, A. Marbeuf, M. Pouchard, P. Hagenmuller, J. Solid State Chem. 1971. 3, 582.
- [2] P. Lacorre, J. B. Torrance, J. Pannetier, A. I. Nazall, P. W. Wang, T. C. Huang, J. Solid State Chem. 1991, 91, 225.

- [3] J. B. Torrance, P. Lacorre, A. I. Nazzal, E. J. Asaldo, C. Niedermayer, *Phys. Rev. B* 1992, 45, 8209.
- [4] J. Gopalakrishnan, G. Colsmann, B. Reuter, J. Solid State Chem. 1977, 22, 145.
- [5] Y. Takeda, R. Kanno, M. Sakano, O. Yamamoto, M. Takano, Y. Bando, H. Akinaga, K. Takita, J. B. Goodenough, Mater. Res. Bull. 1990, 25, 293.
- [6] K. Sreedhar, C. N. Rao, Mater. Res. Bull. 1990, 25, 1235.
- [7] B. W. Arbuckle, K. V. Ramanujachary, Z. Zhang, M. Greenblatt, J. Solid State Chem. 1990, 88, 278.
- [8] Y. Takeda, M. Nishijima, N. Imanishi, R. Kanno, O. Yamamoto, M. Takano, J. Solid State Chem. 1992, 96, 72.
- [9] M. James, J. P. Attfield, J. Mater. Chem. in press.
- [10] S. C. Chen, K. V. Ramanujachary, M. Greenblatt, J. Solid State Chem. 1993, 105, 444.
- [11] M. James, J. P. Attfield, J. Solid State Chem. 1993, 105, 287.
- [12] M. James, J. P. Attfield, J. Rodriguez-Carvajal, Chem. Mater. 1995, 7, 1448.
- [13] M. James, J. P. Attfield, J. Rodriguez-Carvajai, Chem. Ma. [13] M. James, J. P. Attfield, J. Mater. Chem. 1994, 4, 575.
- [14] J. J. Randall, L. Katz, Acta. Crystallogr. 1959, 12, 519.
- [15] J. J. Randall, R. Ward, J. Am. Chem. Soc. 1959, 81, 2629.
- [16] JCPDS Grant-in-Aid-Report; 1908, D. Pfoertsch, Penn State University, Pennsylvania, USA.
- [17] H. Müller-Buschbaum, R. von Schenck, Z. Anorg. Allg. Chem. 1970, 377, 70.
- [18] J. von Wendenburg, A. Henglein, Z. Naturforsch. 1964, 196, 955.
- [19] J. Lee, G. F. Holland, J. Solid State Chem. 1991, 93, 267.
- [20] G. W. Hoffman, J. J. Brown, J. Inorg. Nucl. Chem. 1968, 30, 63.
- [21] R. D. Shannon, Acta Crystallogr. 1976, A32, 751.
- [22] R. Kanno, H. Kubo, Y. Kawamoto, T. Kamiyama, F. Izumi, Y. Takeda, M. Takano, J. Solid State Chem. 1994, 110, 216.
- [23] G. Bergerhoff, O. Schmitz-Dumont, Z. Anorg. Allg. Chem. 1956, 284, 10.
- [24] C. L. McDaniel, J. Am. Ceram. Soc. 1972, 55, 426.
- [25] C. L. McDaniel, J. Solid State Chem. 1972, 4, 275.
- [26] T. N. Nguyen, D. M. Giaquinta, H.-C. zur Loye, Chem. Mater. 1994, 6, 1642.
- [27] A. B. Bykov, S. F. Radaev, E. A. Gerkinai, L. N. Dem'yanets, B. A. Maximov, O. K. Mel'nikov, Kristallografiya 1990, 35, 869.
- [28] A. P. Wilkinson, A. K. Cheetham, W. Kunnman, A. Kvick, Eur. J. Solid State Inorg. Chem. 1991, 28, 453.
- 29] G. Bergerhoff and O. Schmitz-Dumont, Z. Anorg. Allg. Chem. 1956, 284, 10.
- [30] M. D. Joesten, S. Takagi, P. G. Lenhert, Inorg. Chem. 1977, 16, 2680.
- [31] H. Elliott, B. J. Hathaway, R. C. Slade, Inorg. Chem. 1966, 5, 669.
- [32] H. M. Rieveld, J. Appl. Crystallogr. 1969, 2, 65.
- [33] A. C. Larson, R. B. Von Dreele, Los Alamos National Laboratory Report No. LA-UR-86-748, 1987.